

# Intense red upconversion luminescence from $\text{Tm}^{3+}/\text{Yb}^{3+}$ codoped transparent glass ceramic

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$\text{Tm}^{3+}/\text{Yb}^{3+}$  codoped transparent glass ceramic containing  $\beta\text{-PbF}_2$  nanocrystals was successfully prepared. After thermal treatment, emissions from the  $^1G_4$  state of  $\text{Tm}^{3+}$  excited by 980 nm laser were greatly quenched by cross relaxation and the 700 nm luminescence from  $\text{Tm}^{3+}:^3F_{2,3} \rightarrow ^3H_6$  transition was strongly enhanced. A nearly monochromatic red luminescence band was observed. Based on the luminescence decay curves and Judd–Ofelt analysis, the strengthened cross relaxation played an important role in such phenomenon. © 2012 Optical Society of America

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Upconversion (UC) emissions of trivalent rare-earth (RE) ions have attracted a lot of attention for their wide applications in solid-state lasers, sensors, and high-density memories [1–3], and so on. Among the RE ions,  $\text{Tm}^{3+}$  has been extensively investigated for producing UV to IR luminescence [4,5]. Because of the large absorption cross section around 980 nm,  $\text{Tm}^{3+}$  is usually used as the sensitizer. In most previous investigations, the red UC emission (700 nm) from  $\text{Tm}^{3+}:^3F_{2,3} \rightarrow ^3H_6$  is often neglected or cannot be detected due to the low luminescence intensity, but the existence of this emission is theoretically known. Therefore, it is of technological importance and scientific interest to explore this wavelength emission.

The energy level diagram of  $\text{Tm}^{3+}$  and some of the possible cross-relaxation (CR) channels for the  $^1G_4$  state are depicted in Fig. 1. According to the Dexter theory [6], the CR probability is proportional to  $R^{-6}$ , where  $R$  is the distance between RE ions. As a result, such abundant CR channels make the emissions from the  $^1G_4$  state very sensitive to the change of the distance between  $\text{Tm}^{3+}$  ions. Among these CR channels, CR1:  $^1G_4 + ^3H_6 \rightarrow ^3F_{2,3} + ^3F_4$  can easily take place due to the small energy mismatch

(about  $493 \text{ cm}^{-1}$  determined by the absorption spectra in our experiment). Additionally, CR1 can also populate the  $^3F_{2,3}$  levels and further enhance 700 nm UC luminescence. Therefore, it is possible to obtain a bright and monochromatic 700 nm red emission band by the enhancement of CR processes through decreasing the distance between  $\text{Tm}^{3+}$  ions.

We choose the oxyfluoride glass ceramic as the host material in our investigations. Through thermal treatment of the precursor glass, the RE ions can be incorporated into the fluoride nanocrystals, which can effectively reduce the distance between the doping ions. Based on this mechanism, the chemical composition and the sample preparation routine were designed to produce 700 nm UC luminescence.

Transparent glass ceramic was obtained through thermal treatment of the precursor glass. The samples were doped with  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  with relatively high concentration. Under 980 nm laser excitation, our sample yields an intense and nearly monochromatic red luminescence band centering at 700 nm. The mechanisms leading to such interesting phenomenon are investigated by analysis based on the luminescence decay curves and Judd–Ofelt (J-O) theory.

The glass ceramic is prepared with the following starting composition ratios (in mol. %):  $50\text{SiO}_2\text{-}50\text{PbF}_2\text{-}0.5\text{Tm}_2\text{O}_3\text{-}5\text{Yb}_2\text{O}_3$ . The fully mixed chemicals are melted in a covered alumina crucible at  $1000^\circ\text{C}$  for 1 h in a SiC-resistance electric furnace in air. Then, the melt is casted into a preheated stainless steel mold and naturally cooled down to room temperature to form the precursor glass. The precursor glass is heated to  $470^\circ\text{C}$  with a heating rate of  $10 \text{ K/min}$ , and then held for 8 h to form transparent glass ceramic through crystallization. These treatment conditions were set empirically. The resulting samples are cut and polished into the size of  $10 \text{ mm} \times 5 \text{ mm} \times 2 \text{ mm}$  for optical measurements. X-ray diffraction (XRD) patterns were acquired using a Rigaku  $D/\text{max-}\gamma\text{B}$  diffractometer with Cu  $K\alpha$  radiation ( $\lambda = 0.15418 \text{ nm}$ ). The UC luminescence spectra were recorded by a Zolix-SBP300 grating spectrometer equipped with a

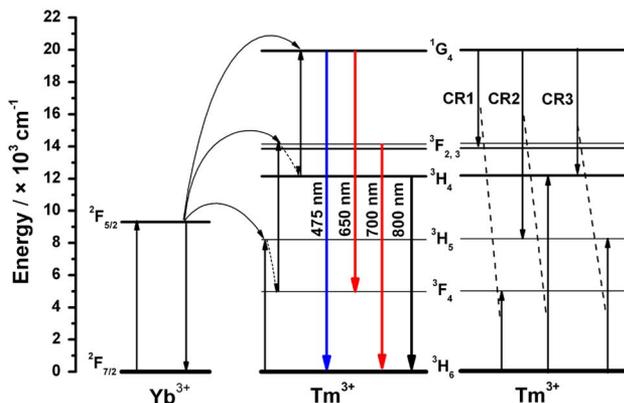


Fig. 1. (Color online) Energy level diagram of  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  and the possible cross-relaxation channels, as well as processes generating blue, red, and near-IR UC emissions.

CR131 photomultiplier tube. A 980 nm diode laser is used as the pump source. The luminescence decay curves were recorded by a Tektronix DPO 4140 oscilloscope and the absorption spectra were measured with a fluorospectrophotometer. All measurements were carried out in ambient pressure and at room temperature.

The XRD patterns of the glass and the glass ceramic are shown in Fig. 2. There are only two humps in the XRD curve for the glass, indicating its amorphous structure. However, after thermal treatment, intense diffraction peaks attributed to  $\beta$ -PbF<sub>2</sub> nanocrystals [7] emerged in the XRD patterns, reflecting the crystallization during thermal treatment. The nanocrystal size is estimated to be about 20 nm in diameter by the Debye–Scherrer formula.

The absorption spectra of samples are shown in Fig. 3. The marked absorption peaks correspond to the transitions from the ground state  $^3H_6$  to the excited states of Tm<sup>3+</sup> and the transition from  $^2F_{7/2}$  to  $^2F_{5/2}$  of Yb<sup>3+</sup>. The inset of Fig. 3 is the enlarged absorption part between 660 and 820 nm. It can be seen that, after crystallization, the  $^3H_6 \rightarrow ^3F_{2,3}$  transition becomes sharper and the absorption maximum for the  $^3H_6 \rightarrow ^3H_4$  transition changes from 791 to 776 nm. These observations indicate that most RE ions have successfully entered into the  $\beta$ -PbF<sub>2</sub> nanocrystals after thermal treatment [8]. Based on the data deduced from the absorption spectra, J-O parameters  $\Omega_t$  ( $t = 2, 4, 6$ ) were determined by the J-O theory [9,10]. Generally speaking, the intensity parameter  $\Omega_2$  is sensitive to the local symmetry environment around RE ions, and it decreases from oxide to fluoride [11]. The calculated value of  $\Omega_2$  (about  $2.76 \times 10^{-20} \text{ cm}^2$ ) in the glass ceramic is smaller than that of the glass (about  $3.40 \times 10^{-20} \text{ cm}^2$ ). These results further verify the incorporation of Tm<sup>3+</sup> and Yb<sup>3+</sup> ions into the  $\beta$ -PbF<sub>2</sub> nanocrystals.

Figure 4 shows the UC emission spectra of both the glass and the glass ceramic excited by a 980 nm laser with its pump power at 350 mW. Emission bands of the glass, which are attributed to the  $^1G_4 \rightarrow ^3H_6$  (475 nm),  $^1G_4 \rightarrow ^3F_4$  (650 nm),  $^3F_{2,3} \rightarrow ^3H_6$  (700 nm), and  $^3H_4 \rightarrow ^3H_6$  (800 nm) transitions of Tm<sup>3+</sup>, are observed. In the visible UC spectrum of the glass, blue emission is dominant, as shown by the inset of Fig. 4(a). For the glass ceramic, the blue and red (650 nm) emissions are greatly quenched so that they become negligible compared to the 700 nm luminescence. On the contrary, the red UC emission from the  $^3F_{2,3} \rightarrow ^3H_6$  transition is enhanced by a factor of 3.5 compared to that of the glass, and the 800 nm near-IR (NIR) emission is also greatly increased. Bright red luminescence can be clearly seen by naked eyes, as presented by the inset of Fig. 4(b). Be-

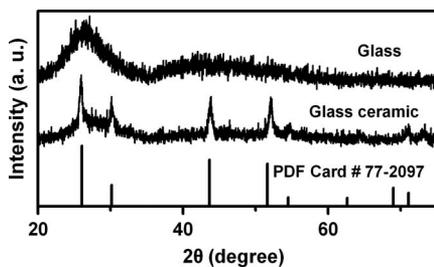


Fig. 2. XRD patterns of the precursor glass and glass ceramic.

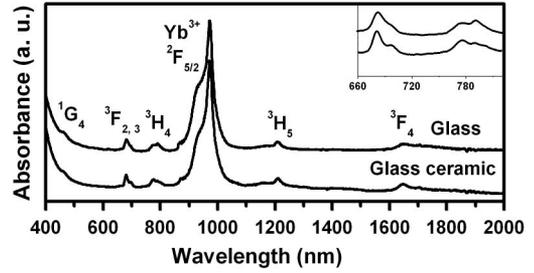


Fig. 3. Absorption spectra for the glass and the glass ceramic; the inset shows the enlarged part between 660 and 820 nm.

cause of the great luminescence quenching of the  $^1G_4$  state, the intensity of 475 and 650 nm emissions keeps much smaller values compared with that of 700 nm luminescence, although the intensity ratio between them can be adjusted by changing the pump power. Therefore, intense 700 nm emission, a quasi-monochromatic red luminescence band in the visible range, has been obtained in the glass ceramic under a 980 nm laser excitation.

The weakening of emissions from the Tm<sup>3+</sup>: $^1G_4$  state when the glass becomes glass ceramic probably originate from the enhanced cross relaxations by crystallization. To confirm this mechanism, the cross-relaxation rate was calculated by the following equation:

$$1/\tau = 1/\tau_{\text{rad}} + W_{\text{ph}} + W_{\text{cr}}, \quad (1)$$

where  $\tau$  is the luminescence decay time;  $\tau_{\text{rad}}$  is the radiation lifetime,  $W_{\text{ph}}$  is the phonon-assisted nonradiation relaxation rate,  $W_{\text{cr}}$  is the cross-relaxation rate. The luminescence decay curves for the Tm<sup>3+</sup>: $^1G_4$  state in the glass and glass ceramic are shown in Fig. 5(a). The values of  $\tau$  are 150 and 45  $\mu\text{s}$  for the  $^1G_4$  state in the glass and in the glass ceramic, respectively.  $W_{\text{ph}}$  can be ignored due to the large energy gap from  $^1G_4$  to the next level. The parameter  $\tau_{\text{rad}}$  can be calculated based on the J-O theory [12]:

$$\tau_{\text{rad}} = \frac{1}{A_{\text{ed}} + A_{\text{md}}}, \quad (2)$$

$$A_{\text{ed}} = \frac{64\pi^4 e^2}{3h(2J+1)\lambda^3} \times \frac{n(n^2+2)^2}{9} \sum_{t=2,4,6} \Omega_t |\langle \gamma J \| U^{(t)} \| \gamma' J' \rangle|^2, \quad (3)$$

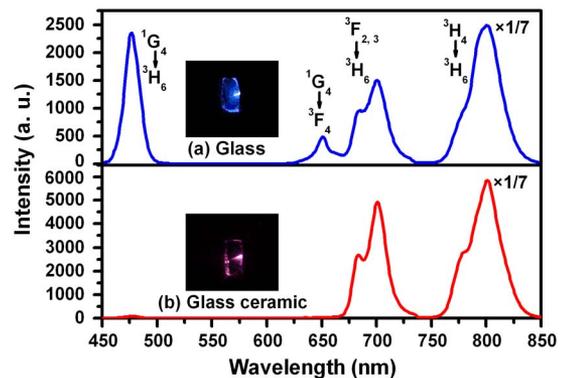


Fig. 4. (Color online) Upconversion emission spectra of (a) the precursor glass and (b) the glass ceramic under 980 nm laser excitation; the insets of (a) and (b) show photographs of luminescence from the glass and the glass ceramic.

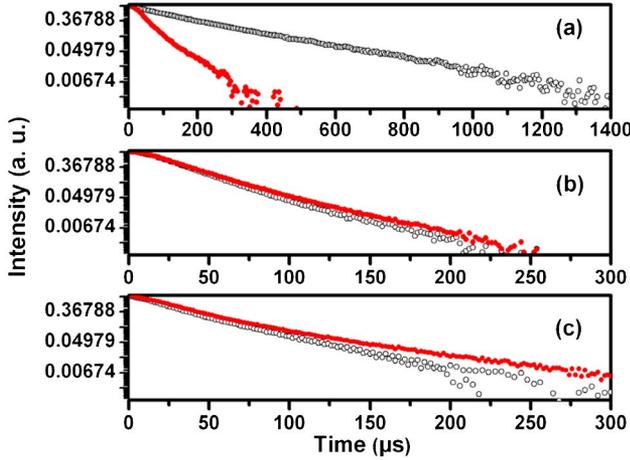


Fig. 5. (Color online) Luminescence decay curves for the (a)  $^1G_4$ , (b)  $^2F_{2,3}$ , and (c)  $^3H_4$  states in the precursor glass (open black symbols) and glass ceramic (solid red symbols).

$$A_{\text{md}} = \frac{4\pi^2 e^2 h n^3}{3m^2 c^2 (2J+1)\lambda^3} |\langle \gamma J \| L + 2S \| \gamma' J' \rangle|^2. \quad (4)$$

The meaning of the parameters in Eqs. (2), (3), and (4) can be consulted in [12]. Finally, the calculated  $W_{\text{cr}}$  for the  $^1G_4$  state in the glass ceramic is about  $19.999 \times 10^3 \text{ s}^{-1}$ , which is about 4.7 times the value for the glass ( $4.286 \times 10^3 \text{ s}^{-1}$ ). This result confirms the enhancement of CR processes after crystallization. Furthermore, it is worth pointing out that the measured luminescence lifetime for  $^3F_{2,3}$  states is hardly changed (from 23 to 24  $\mu\text{s}$ ) after thermal treatment, as shown in Fig. 5(b). This means that the enhancement of 700 nm UC luminescence is mainly due to the increased population of the  $^3F_{2,3}$  states, which is attributed to several reasons. First, the increased CR1:  $^1G_4 + ^3H_6 \rightarrow ^3F_{2,3} + ^3F_4$  plays an important role in weakening emissions from the  $^1G_4$  state and enhancing the population in the  $^3F_{2,3}$  states. Second, the much-enhanced energy transfer (ET) from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$  in the glass ceramic also contributes. After crystallization, the distance between  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  is reduced because of the incorporation of RE ions into the fluoride nanocrystals. Based on Dexter theory [13], a shorter distance between sensitizer and acceptor is of benefit for the enhancement of ET efficiency. As a result, the efficiency of the excitation state absorption process  $^3F_4 \rightarrow ^3F_{2,3}$  via ET from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$  is increased, and contributes to the population in the  $^3F_{2,3}$  states. Third, the enhanced population of the  $^3F_4$  state after crystallization also contributes. Through the CR

processes (CR1, CR2, CR3, and  $^3H_4 + ^3H_6 \rightarrow ^3F_4 + ^3F_4$ ), more  $\text{Tm}^{3+}$  ions can be populated in the  $^3F_4$  state. Some of these ions come back to the ground state by radiative or nonradiative transitions, and some are being pumped to the  $^3F_{2,3}$  states via ET from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$ , which finally increases the population of  $^3F_{2,3}$  states. Meanwhile, through the  $\text{Tm}^{3+}: ^3H_4$  state luminescence decay curves in the precursor glass and glass ceramic shown in Fig. 5(c) (changing from 44 to 48  $\mu\text{s}$ ), it is clearly seen that these CR processes also give important contribution to the enhancement of 800 nm NIR emission.

In conclusion, we have demonstrated that intense and nearly monochromatic red UC luminescence band centering at 700 nm can be produced via crystallization of  $\text{Tm}^{3+}/\text{Yb}^{3+}$  codoped  $\text{SiO}_2\text{-PbF}_2$  oxyfluoride glass. The UC emissions from the  $\text{Tm}^{3+}: ^1G_4$  state are greatly quenched by cross relaxation, which is confirmed by theoretical analysis based on the luminescence decay curves and J-O theory. The increased CR processes combine with higher efficiency of ET from  $\text{Yb}^{3+}$  to  $\text{Tm}^{3+}$ , which are induced by the RE incorporation into the  $\beta\text{-PbF}_2$  nanocrystals, greatly enhance the population of  $^3F_{2,3}$  states, and, consequently, increase 700 nm red UC luminescence. This interesting characteristic of  $\text{Tm}^{3+}/\text{Yb}^{3+}$  codoped  $\beta\text{-PbF}_2$  transparent glass ceramic may be used to design devices to meet specific color demands, such as in the fields of optics and solid-state lasers.

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